

Universal Quantum Computation with Gapped Boundaries

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C, Cheng, Wang, *Phys. Rev. Lett.* **119**, 170504 (2017).

Introduction: Quantum Computation

- Quantum computing:
 - Feynman (1982): Use a quantum-mechanical computer to simulate quantum physics (classically intractable)
 - Encode information in qubits, gates = unitary operations
 - Applications: Factoring/breaking RSA (Shor, 1994), quantum machine learning, ...

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 - Feynman (1982): Use a quantum-mechanical computer to simulate quantum physics (classically intractable)
 - Encode information in qubits, gates = unitary operations
 - Applications: Factoring/breaking RSA (Shor, 1994), quantum machine learning, ...
- Major challenge: local decoherence of qubits

Introduction: Topological Quantum Computation

- **Topological** quantum computing (TQC) (Kitaev, 1997; Freedman et al., 2003):
 - Encode information in *topological* degrees of freedom
 - Perform *topologically protected* operations



Topological Quantum Computation

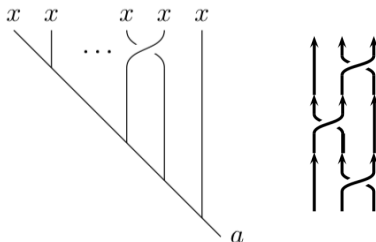
- Traditional realization of TQC: *anyons* in *topological phases of matter* (e.g. Fractional Quantum Hall - FQH)
 - Elementary (quasi-)particles in 2 dimensions s.t. $|\psi_1\psi_2\rangle = e^{i\phi}|\psi_2\psi_1\rangle$

Topological Quantum Computation

- Traditional realization of TQC: *anyons* in *topological phases of matter* (e.g. Fractional Quantum Hall - FQH)
 - Elementary (quasi-)particles in 2 dimensions s.t. $|\psi_1\psi_2\rangle = e^{i\phi}|\psi_2\psi_1\rangle$
- Qubit encoding: degeneracy arising from the *fusion rules*
 - e.g. Toric code: $e^i m^j \otimes e^k m^l = e^{(i+k)(\bmod 2)} m^{(j+l)(\bmod 2)}$ ($i, j = 0, 1$)
 - e.g. Ising: $\sigma \otimes \sigma = 1 \oplus \psi$, $\psi \oplus \psi = 1$
 - e.g. Fibonacci: $\tau \otimes \tau = 1 \oplus \tau$

Topological Quantum Computation

- Topologically protected operations: Braiding of anyons
 - Move one anyon around another \rightarrow pick up phase (due to Aharonov-Bohm)



Figures: (1) Z. Wang, *Topological Quantum Computation*.

(2) C. Nayak et al., *Non-abelian anyons and topological quantum computation*.

Problem

- Problem: Abelian anyons have no degeneracy, so no computation power \rightarrow need non-abelian anyons for TQC (e.g. $\nu = 5/2, 12/5$)

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- These are difficult to realize, existence is still uncertain
- Question: Given a top. phase that supports only abelian anyons, is it possible “engineer” other non-abelian objects?
 - Answer: Yes! We consider boundaries of the topological phase \rightarrow *gapped boundaries*
 - We'll even get a universal gate set from gapped boundaries of an abelian phase (specifically, bilayer $\nu = 1/3$ FQH)

- Framework of TQC
- Introduce gapped boundaries and their framework
- Gapped boundaries for TQC
- Universal TQC with gapped boundaries in bilayer $\nu = 1/3$ FQH
- Summary and Outlook

- **Framework of TQC**
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Framework of TQC

Formally, an *anyon model* \mathcal{B} consists of the following data:

- Set of anyon types/labels: $\{a, b, c \dots\}$, one of which should represent the vacuum **1**
 - Each anyon type has a topological twist $\theta_i \in U(1)$:

$$\begin{array}{c} \curvearrowright \\ i \end{array} = \theta_i \begin{array}{c} | \\ i \end{array}$$

¹More general cases exist, but are not used in this talk.

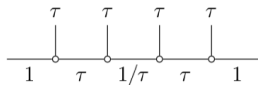
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- Set of anyon types/labels: $\{a, b, c, \dots\}$, one of which should represent the vacuum $\mathbf{1}$
 - Each anyon type has a topological twist $\theta_i \in U(1)$:

$$\text{loop}_i = \theta_i \mid_i$$

- For each pair of anyon types, a set of fusion rules: $a \otimes b = \bigoplus_c N_{ab}^c c$.
 - The fusion space of $a \otimes b$ is a vector space V_{ab} with basis V_{ab}^c .¹
 - The fusion space of $a_1 \otimes a_2 \otimes \dots \otimes a_n$ to b is a vector space $V_{a_1 a_2 \dots a_n}^b$ with basis given by anyon labels in intermediate segments, e.g.



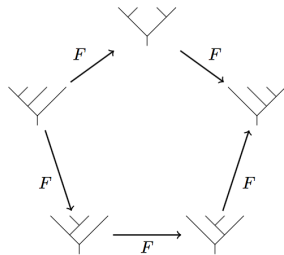
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Framework of TQC

Formally, an *anyon model* \mathcal{B} consists of the following data: [Cont'd]

- Associativity: for each (a, b, c, d) , a set of (unitary) linear transformations $\{F_{d;ef}^{abc} : V_d^{abc} \rightarrow V_d^{abc}\}$ satisfying “pentagons”

$$\begin{array}{c} i \quad j \quad k \\ \diagdown \quad \diagup \quad \diagup \\ m \quad \quad \quad \\ \diagup \quad \diagdown \\ l \end{array} = \sum_n F_{l;nm}^{ijk} \begin{array}{c} i \quad j \quad k \\ \diagdown \quad \diagup \quad \diagup \\ \quad \quad n \quad \\ \diagup \quad \diagdown \\ l \end{array}$$

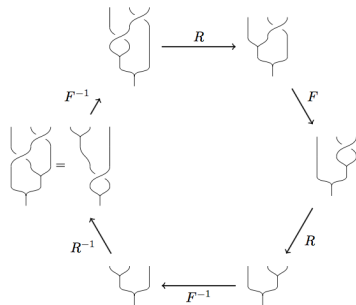


Framework of TQC

Formally, an *anyon model* \mathcal{B} consists of the following data: [Cont'd]

- (Non-degenerate) braiding: For each (a, b, c) s.t. $N_{a,b}^c \neq 0$, a set of phases¹ $R_{ab}^c \in U(1)$ compatible with associativity (“hexagons”):

$$\begin{array}{c} a \quad b \\ \diagdown \quad \diagup \\ \text{---} \text{---} \\ \diagup \quad \diagdown \\ c \end{array} = R_c^{ab} \begin{array}{c} a \quad b \\ \diagdown \quad \diagup \\ \text{---} \text{---} \\ \diagup \quad \diagdown \\ c \end{array}$$



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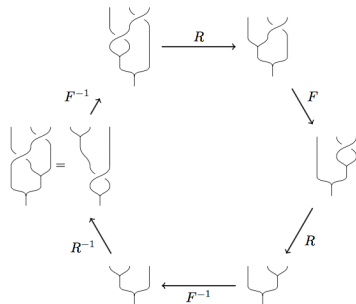
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Mathematically, this is captured by a (unitary) *modular tensor category* (UMTC)



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Framework of TQC

- In this framework, a 2D *topological phase of matter* is an equivalence class of gapped Hamiltonians $\mathcal{H} = \{H\}$ whose low-energy excitations form the same anyon model \mathcal{B}
- Examples (on the lattice) include Kitaev's quantum double models, Levin-Wen string-net models, ...

Framework of TQC: Example

- Physical system: Bilayer FQH system, $1/3$ Laughlin state of opposite chirality in each layer
- Equivalent to \mathbb{Z}_3 toric code (Kitaev, 2003)
- Anyon types: $e^a m^b$, $a, b = 0, 1, 2$
- Twist: $\theta(e^a m^b) = \omega^{ab}$ where $\omega = e^{2\pi i/3}$
- Fusion rules: $e^a m^b \otimes e^c m^d \rightarrow e^{(a+c)(\bmod 3)} m^{(b+d)(\bmod 3)}$
- F symbols all trivial (0 or 1)
- $R_{e^a m^b, e^c m^d} = e^{2\pi i bc/3}$
- UMTC: $SU(3)_1 \times \overline{SU(3)_1} \cong \mathfrak{D}(\mathbb{Z}_3) = \mathcal{Z}(\text{Rep}(\mathbb{Z}_3)) = \mathcal{Z}(\text{Vec}_{\mathbb{Z}_3})$

- Framework of TQC
- **Introduce gapped boundaries and their framework**
- Gapped boundaries for TQC
- Universal gate set with gapped boundaries in bilayer $\nu = 1/3$ FQH
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Gapped Boundaries: Framework

- A *gapped boundary* is an equivalence class of gapped local (commuting) extensions of $H \in \mathcal{H}$ to the boundary

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- In the anyon model: Collection of bulk bosonic ($\theta = 1$) anyons which condense to vacuum on the boundary (think Bose condensation)
 - All other bulk anyons condense to confined “boundary excitations”
 $\alpha, \beta, \gamma \dots$
- Mathematically, Lagrangian algebra $\mathcal{A} \in \mathcal{B}$

Gapped Boundaries: Framework

More rigorously, gapped boundaries come with M symbols (like F symbols for the bulk):

$$\begin{array}{c} a \quad b \\ | \quad | \\ \hline \end{array} = \sum_c M_c^{ab} \begin{array}{c} a \quad b \\ \diagdown \quad \diagup \\ c \\ | \\ \hline \end{array}$$

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(More generally, we also define these with boundary excitations, but that is unnecessary for this talk.)

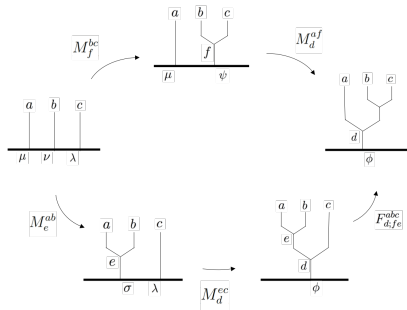
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M symbols must be compatible with F symbols (“mixed pentagons”):



Gapped Boundaries: Example

- Physical system: Bilayer $\nu = 1/3$ FQH, equiv. to \mathbb{Z}_3 toric code
- Anyon types: $e^a m^b$, $a, b = 0, 1, 2$

Gapped Boundaries: Example

- Physical system: Bilayer $\nu = 1/3$ FQH, equiv. to \mathbb{Z}_3 toric code
- Anyon types: $e^a m^b$, $a, b = 0, 1, 2$
- Two gapped boundary types:
 - Electric charge condensate: $\mathcal{A}_1 = 1 \oplus e \oplus e^2$
 - Magnetic flux condensate: $\mathcal{A}_2 = 1 \oplus m \oplus m^2$

Gapped Boundaries: Example

- Physical system: Bilayer $\nu = 1/3$ FQH, equiv. \mathbb{Z}_3 toric code
- We will work mainly with $\mathcal{A}_1 = 1 \oplus e \oplus e^2$:
 - Algebraically, $\mathcal{A}_1, \mathcal{A}_2$ are equivalent by electric-magnetic duality
 - Easier to work with charge condensate - read-out can be done by measuring electric charge (Barkeshli, 2016)
 - It is interesting to consider *both* \mathcal{A}_1 and \mathcal{A}_2 at the same time - we do this in a separate paper³, will briefly mention in our Outlook

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- M symbols for this theory are all 0 or 1

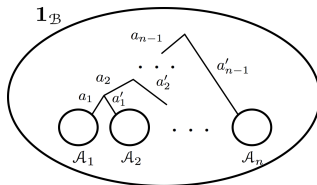
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Gapped Boundaries for TQC

Qudit encoding:

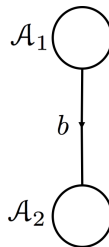
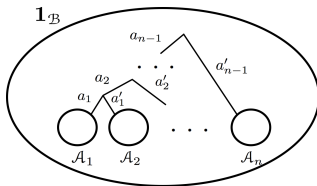
- Gapped boundaries give rise to a natural ground state degeneracy: n gapped boundaries on a plane, with total charge vacuum



Gapped Boundaries for TQC

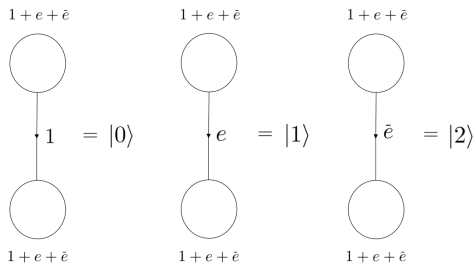
Qudit encoding:

- Gapped boundaries give rise to a natural ground state degeneracy: n gapped boundaries on a plane, with total charge vacuum
- For qudit encoding: use $n = 2$



Gapped Boundaries for TQC: Example

For our bilayer $\nu = 1/3$ FQH system, we have:



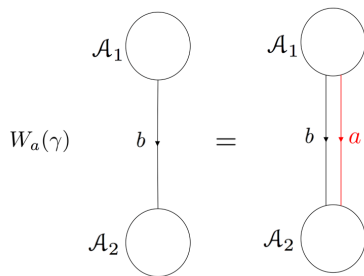
Gapped Boundaries for TQC

Topologically protected operations:

- Tunnel- a operations
- Loop- a operations
- Braiding gapped boundaries
- Topological charge measurement*

Tunnel-a Operations

Starting from state $|b\rangle$, tunnel an a anyon from \mathcal{A}_1 to \mathcal{A}_2 :



Tunnel-a Operations

Compute using M -symbols:

$$\begin{array}{c} \mathcal{A}_1 \\ \downarrow \\ b \rightarrow \text{red line} \leftarrow a \\ \downarrow \\ \mathcal{A}_2 \end{array} = \sum_c M_c^{ab}(\mathcal{A}_2) [M_c^{ab}(\mathcal{A}_1)]^\dagger
 \begin{array}{c} \mathcal{A}_1 \\ \downarrow c \\ \begin{array}{c} \downarrow \\ b \rightarrow \quad \leftarrow a \\ \downarrow \\ c \end{array} \\ \downarrow c \\ \mathcal{A}_2 \end{array}$$

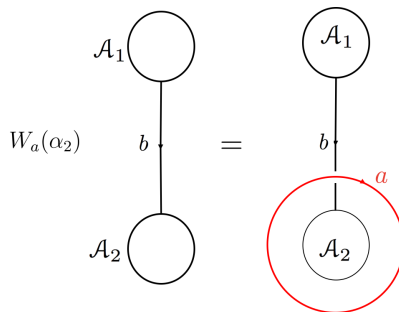
Tunnel-a Operations

Result:

$$W_a(\gamma) \begin{array}{c} \mathcal{A}_1 \\ \downarrow b \\ \mathcal{A}_2 \end{array} = \sum_c M_c^{ab}(\mathcal{A}_1) [M_c^{ab}]^\dagger(\mathcal{A}_2) \sqrt{\frac{d_a d_b}{d_c}} \begin{array}{c} \boxed{\mathcal{A}_1} \\ \downarrow c \\ \boxed{\mathcal{A}_2} \end{array} \quad (1)$$

Loop- a Operations

Starting from state $|b\rangle$, loop an a anyon around one of the boundaries:



Loop-a Operations

Similar computation methods lead to the formula:

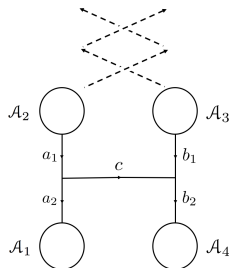
$$W_a(\alpha_2) \begin{array}{c} \text{---} \mathcal{A}_1 \\ | \\ b \cdot \\ | \\ \text{---} \mathcal{A}_2 \end{array} = s_{ab} \begin{array}{c} \text{---} \mathcal{A}_1 \\ | \\ b \cdot \\ | \\ \text{---} \mathcal{A}_2 \end{array}$$

where $s_{ab} = \tilde{s}_{ab}/d_b$ is given by the modular S matrix of the theory:

$$\tilde{s}_{ij} = \begin{array}{c} \text{---} \bigcirc \bigcirc \text{---} \\ i \quad j \end{array}$$

Braiding Gapped Boundaries

Braid gapped boundaries around each other:



(Mathematically, this gives a representation of the (spherical) $2n$ -strand pure braid group.)

Braiding Gapped Boundaries

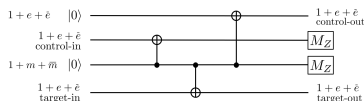
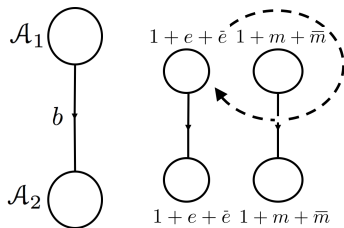
Simplify with (bulk) R and F moves to get:

$$\sigma_2^2 \begin{array}{c} \mathcal{A}_2 \quad \mathcal{A}_3 \\ \begin{array}{c} a_1 \\ a_2 \end{array} \quad \begin{array}{c} c \\ b_1 \\ b_2 \end{array} \\ \mathcal{A}_1 \quad \mathcal{A}_4 \end{array} = \sum_{c, c'} F_{b_2; c' c}^{a_2 a_1 b_1} R_{c'}^{b_1 a_1} R_c^{a_1 b_1} (F_{b_2}^{a_2 a_1 b_1})_{c'' c'}^{-1} \begin{array}{c} \mathcal{A}_2 \quad \mathcal{A}_3 \\ \begin{array}{c} a_1 \\ a_2 \end{array} \quad \begin{array}{c} c'' \\ b_1 \\ b_2 \end{array} \\ \mathcal{A}_1 \quad \mathcal{A}_4 \end{array} \quad (2)$$

Gapped Boundaries for TQC: Example

For our bilayer $\nu = 1/3$ FQH case: ($\mathcal{A}_1 = \mathcal{A}_2 = 1 \oplus e \oplus e^2$)

- Tunnel an e anyon from \mathcal{A}_1 to \mathcal{A}_2 : $W_a(\gamma)|b\rangle = |a \otimes b\rangle$
 $\rightarrow W_e(\gamma) = \sigma_3^x$, where $\sigma_3^x|i\rangle = |(i+1)(\text{mod } 3)\rangle$
- Loop an m anyon around \mathcal{A}_2 : $W_m(\alpha_2)|e^j\rangle = \omega^j|e^j\rangle$
 $\rightarrow W_m(\alpha_2) = \sigma_3^z$
- Braid gapped boundaries: get $\wedge \sigma_3^z$

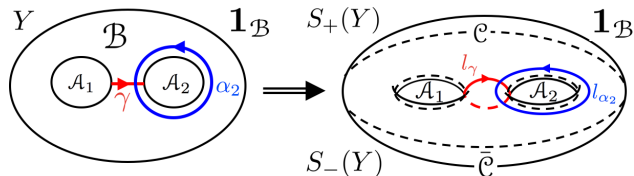


Topological Charge Projection/Measurement

- Motivation:
 - Property F conjecture (Naidu and Rowell, 2011): Braiding alone cannot be universal for TQC for most physically plausible systems

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 - Property F conjecture (Naidu and Rowell, 2011): Braidings alone cannot be universal for TQC for most physically plausible systems
- Topological charge projection (TCP) (Barkeshli and Freedman, 2016):
 - Doubled theories: Wilson line lifts to a loop \rightarrow measure topological charge through the loop



- Resulting projection operator: $(\mathcal{O}_x(\beta) = W_{x\bar{x}}(\gamma_i) \text{ or } W_x(\alpha_i))$

$$P_{\beta}^{(a)} = \sum_{x \in \mathcal{C}} S_{0a} S_{xa}^* \mathcal{O}_x(\beta). \quad (3)$$

Topological Charge Projection/Measurement

- Topological charge projection (TCP): [Cont'd]
 - Given an anyon theory \mathcal{C} , its \mathcal{S}, \mathcal{T} matrices

$$\mathcal{S} = \left\{ \tilde{s}_{ij} = \begin{array}{c} \text{diagram of two overlapping circles} \\ i \quad j \end{array} \right\}, \quad \mathcal{T} = \text{diag}(\theta_i)$$

give mapping class group representations $V_{\mathcal{C}}(Y)$ for surfaces Y .

- Barkeshli and Freedman showed that topological charge projections generate all matrices in $V_{\mathcal{C}}(Y)$

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give mapping class group representations $V_{\mathcal{C}}(Y)$ for surfaces Y .

- Barkeshli and Freedman showed that topological charge projections generate all matrices in $V_{\mathcal{C}}(Y)$
- General topological charge measurements (TCMs):
 - Projection operators $P_{\beta}^{(a)} = \sum_{x \in \mathcal{C}} S_{0a} S_{xa}^* \mathcal{O}_x(\beta) \rightarrow$ topological charge measurements perform the *complement* of $P_{\beta}^{(a)1}$
 - Not always physical, but special cases are *symmetry protected* – we examine this

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Universal Gate Set

Universal (*metaplectic*) gate set for the qutrit model (Cui and Wang, 2015):

- 1 The single-qutrit Hadamard gate H_3 , defined as $H_3|j\rangle = \frac{1}{\sqrt{3}} \sum_{i=0}^2 \omega^{ij}|i\rangle$, $j = 0, 1, 2$, $\omega = e^{2\pi i/3}$
- 2 The two-qutrit entangling gate SUM_3 , defined as $\text{SUM}_3|i\rangle|j\rangle = |i\rangle|(i+j) \bmod 3\rangle$, $i, j = 0, 1, 2$.
- 3 The single-qutrit generalized phase gate $Q_3 = \text{diag}(1, 1, \omega)$.
- 4 Any nontrivial single-qutrit classical (i.e. Clifford) gate not equal to H_3^2 .
- 5 A projection M of a state in the qutrit space \mathbb{C}^3 to $\text{Span}\{|0\rangle\}$ and its orthogonal complement $\text{Span}\{|1\rangle, |2\rangle\}$, so that the resulting state is coherent if projected into $\text{Span}\{|1\rangle, |2\rangle\}$.

Universal Gate Set - Bilayer $\nu = 1/3$ FQH

Universal (*metaplectic*) gate set for the qutrit model (Cui and Wang, 2015):

- ① $H_3 = \mathcal{S}$ for $\mathcal{C} = \text{SU}(3)_1$ (single layer $\nu = 1/3$ FQH), so it can be implemented by TCP.

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- ② $\wedge\sigma_3^z$ can be implemented by gapped boundary braiding. Conjugate second qutrit by H_3 to get SUM_3 .

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- ③ TCP can implement $\text{diag}(1, \omega, \omega)$ (Dehn twist of $\text{SU}(3)_1$). Follow by σ_3^x for Q_3 .

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Universal (*metaplectic*) gate set for the qutrit model (Cui and Wang, 2015):

- ① $H_3 = \mathcal{S}$ for $\mathcal{C} = \text{SU}(3)_1$ (single layer $\nu = 1/3$ FQH), so it can be implemented by TCP.
- ② $\wedge \sigma_3^z$ can be implemented by gapped boundary braiding. Conjugate second qutrit by H_3 to get SUM_3 .
- ③ TCP can implement $\text{diag}(1, \omega, \omega)$ (Dehn twist of $\text{SU}(3)_1$). Follow by σ_3^x for Q_3 .
- ④ Any nontrivial single-qutrit classical (i.e. Clifford) gate not equal to H_3^2 - we have a Pauli-X from tunneling e .
- ⑤ Projective measurement - we use the TCM which is the complement of

$$P_\gamma^{(1)} = \frac{1}{3} \begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix}. \quad (4)$$

Conjugating $1 - P_\gamma^{(1)}$ with the Hadamard gives the result.

Universal Gate Set - Bilayer $\nu = 1/3$ FQH

To get the projective measurement, we introduce a *symmetry-protected* topological charge measurement:

- Want to tune system s.t. quasiparticle tunneling along γ is enhanced
→ implement $H' = -tW_\gamma(e) + \text{h.c.}$
 - $t = (\text{complex})$ tunneling amplitude, $W_\gamma(e) = \text{tunnel-e operator}$

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- Implementing $M \leftrightarrow$ ground state of H' is doubly degenerate for $|e\rangle, |e^2\rangle \leftrightarrow t$ is real (beyond topological protection)
- Physically, could realize in fractional quantum spin Hall state – quantum spin Hall + time-reversal symmetry (exchange two layers)
 - Topologically equiv. to $\nu = 1/3$ Laughlin, $e =$ bound state of spin up/down quasiholes
 - e is time-reversal invariant → tunneling amplitude of e must be real
→ symmetry-protected TCM

- Framework of TQC
- Introduce gapped boundaries and their framework
- Gapped boundaries for TQC
- Universal gate set with gapped boundaries in bilayer $\nu = 1/3$ FQH
- **Summary and Outlook**

Summary

- Decoherence is a major challenge to quantum computing → **topological quantum computing** (TQC)
- TQC with anyons requires non-abelian topological phases (difficult to implement) → **engineer non-abelian objects** (e.g. **gapped boundaries**) from abelian phases
- We can get a **universal quantum computing gate set** from a purely abelian theory (bilayer $\nu = 1/3$ FQH), which is trivial for anyonic TQC
 - Topologically protected qudit encoding and Clifford gates
 - Symmetry-protected implementation for non-Clifford projection

- Practical implementation of the symmetry-protected TCM
- More thorough study of symmetry-protected quantum computation
 - Amount of protection offered and computation power
- Other routes to engineer non-abelian objects
 - Boundary defects/parafermion zero modes from gapped boundaries of $\nu = 1/3$ FQH (Lindner et al., 2014)
 - Genons and symmetry defects (Barkeshli et al., 2014; C, Cheng, Wang, 2017)
 - How would these look when combined with gapped boundaries?

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Thanks!